



# Article Synthesized Transmitting Coil for Magnetic Focusing of Pulsed Eddy Current for Downhole Casing Inspection

Ling Yang <sup>1</sup>, Changzan Liu <sup>1</sup>, Jingxin Dang <sup>2</sup>, Yang Zhao <sup>1</sup>, Bo Dang <sup>1,\*</sup> and Ruirong Dang <sup>1,\*</sup>

- <sup>1</sup> Shaanxi Key Laboratory of Measurement and Control Technology for Oil and Gas Wells, Xi'an Shiyou University, Xi'an 710065, China; lingyang2915@163.com (L.Y.); liuchang-zan@mail.nwpu.edu.cn (C.L.); 18892892786@163.com (Y.Z.)
- <sup>2</sup> School of Resources and Environment, University of Electronic Science and Technology of China, Chengdu 611731, China; 15029276190@163.com
- \* Correspondence: dangbo@xsyu.edu.cn (B.D.); dangrr@xsyu.edu.cn (R.D.); Tel.: +86-029-8838-2648 (B.D. & R.D.)

**Abstract:** Pulsed eddy current (PEC) is a widely utilized technology for the nondestructive inspection of industrial tubes and pipes due to its rapid and accurate results. To improve the longitudinal resolution of PEC, multiple transmitting coils (MTCs) are used to realize magnetic focusing. However, this approach is difficult to apply to narrow downhole environments because of the complex transmitting array and electrical circuits. To address this issue, we present a synthesized transmitting coil (STC) that combines MTCs into a single coil with multiple connected sections using different winding directions and number of turns to adjust the magnetic field distribution. A theoretical derivation was presented for the analysis and interpretation of the magnetic field, and a figure of merit (FoM) was constructed to optimize the STC parameters. Numerical simulations and experimental results showed good agreement with the simulation results.

**Keywords:** magnetic focusing; synthesized transmitting coil; downhole; pulsed eddy current; casing inspection

# 1. Introduction

Pulsed eddy current (PEC) is a widely used technology for inspecting borehole casings due to its rapid and accurate acquisition of data over a broad frequency range and accessibility to targets [1,2]. The basic PEC probe comprises transmitting coils to excite the magnetic field in the conductive tube being tested, as well as a receiving sensor that picks up the magnetic field from the eddy currents induced by the abrupt cutoff of emissions [3,4]. Many researchers have focused on designing different PEC probes to improve the performance of inspections [5–8]. On the receiver side, different receiving sensors have been utilized to improve the signal-to-noise ratio (SNR), including the coil sensor [5], the Hall sensor [6], and the magnetic resistance sensor [7]. On the transmitter side, research has focused on optimizing the shape and geometric parameters of the sensor design as performance is affected by the magnetic field distribution [8]. Fu et al. [9] used a longitudinal probe and transverse probe to generate magnetic fields in different directions, and the experimental results showed that the orientation of the magnetic field distribution had a significant influence on the detection sensitivity for cracks in the casing. Wu et al. [10] analyzed the field and impedance of a rotary coil in a conductive tube considering the induced current. Zhou et al. [11] analyzed the magnetic flux density distribution of different types of probes to optimize the performance of the transmitting coils. Dang et al. [12] analyzed the effect of the probe size on the eddy current field at different times, and their results indicated that the longitudinal spatial resolution could be improved by reducing the longitudinal distribution range of the magnetic field.



Citation: Yang, L.; Liu, C.; Dang, J.; Zhao, Y.; Dang, B.; Dang, R. Synthesized Transmitting Coil for Magnetic Focusing of Pulsed Eddy Current for Downhole Casing Inspection. *Appl. Sci.* 2022, *12*, 7695. https://doi.org/10.3390/ app12157695

Academic Editors: Cesar Levy and Dwayne McDaniel

Received: 12 July 2022 Accepted: 28 July 2022 Published: 30 July 2022

**Publisher's Note:** MDPI stays neutral with regard to jurisdictional claims in published maps and institutional affiliations.



**Copyright:** © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). In recent years, magnetic focusing (MF) has attracted attention as an effective method of controlling the magnetic field distribution. A common approach to realizing MF is based on magnetic shielding theory [13,14]. Yang et al. [13] proposed a multilayer MF sensor framework with a driver coil, a pickup coil, and a three-layer shield with different shielding materials to enhance the detectability of ferromagnetic materials at greater depths. Tsukada et al. [14] applied high-temperature superconducting plates together as a magnetic shield to

concentrate the magnetic flux inside of the slit between the plates to improve the sensitivity and noise characteristics of a tunnel magnetoresistive sensor. Another approach to realizing MF is based on beamforming or spatial signal processing, which is similar to radar, sonar, and wireless communications [15–19]. This approach uses multiple transmitting coils (MTCs) to control the vector of the magnetic field. Kim et al. [15,16] proposed a method for calculating the current distribution, which enables active sidelobe suppression of MF while maintaining the resolution.

However, while MF is suitable for the application of PEC at the surface, the high temperature, high pressure, and narrow space of downhole environments significantly limit the applicability of MF to PEC in metal pipelines [20–22]. With regard to magnetic shielding, the effect of MF largely depends on the geometric design of the shield [13]. To ensure trafficability [20], a downhole PEC instrument needs to be much smaller than the inner diameter of the casing. This results in some distance between the instrument and the casing, which can cause a poor focusing effect. Meanwhile, MF using MTCs requires different transmitting currents and corresponding electrical circuits, which both increase the hardware cost, and are difficult to fit in the narrow underground space [17,21].

To solve this problem, we propose a synthesized transmitting coil (STC) for PEC downhole casing inspection that combines MTCs into a single coil with multiple connected sections using different winding directions and number of turns to adjust the magnetic field distribution. A figure of merit (FoM) was developed to determine the optimal STC parameters, and simulations and experiments were performed to verify the proposed STC. Section 2 presents the working principle and model of a PEC system for downhole casing inspection. Section 3 presents the proposed STC and FoM used to evaluate its performance. The simulations and results of the experiments are discussed in Section 4. Finally, the paper is concluded in Section 5.

#### 2. Working Principle and Model of the PEC System

Figure 1 shows a PEC system used for borehole casing inspection. The PEC system is suspended on a cable and is controlled by a winch during inspection. Each measurement period corresponds to a depth along the borehole axis. The PEC system is based on the principle of electromagnetic induction. When a direct current is applied to the transmitting coil, a primary magnetic field is formed in a cylindrical layered structure comprising air, the metal casing, cement, and the formation. The current is then turned off instantaneously, and the abrupt change in the magnetic field induces eddy currents in the cylindrical medium. In this paper, the turn-on, turn-off, and on-time of the input current were set to 50  $\mu$ s, 30  $\mu$ s, and 200 ms, respectively. This produces a secondary magnetic field that can be captured by the receiving sensor. The average thickness over the distribution range of the transmitting magnetic field is indicated by the decay rate of the electromagnetic signal.

The PEC system can be modeled as a probe comprising the transmitting coil and receiving sensor that are located at the center of the borehole casing. The relative magnetic permeability, relative dielectric constant, electrical conductivity, and radius of the *j*th layer can be defined as  $\mu_j$ ,  $\varepsilon_j$ ,  $\sigma_j$ , and  $r_j$ , respectively, with j = 1, 2, ..., J. Assuming that the transmitting current is *I*, then the magnetic field at the position of the borehole axis *z* with radius *r* excited by one turn of the coil with the position  $z_t$  can be calculated as follows [23]:



Figure 1. PEC system for borehole casing inspection.

$$H_1(z_t,\omega,r,z) = -\frac{Ir_1}{\pi} \int_0^\infty x K_0(x_j r) I_1(x_j r_1) \cos[\lambda(z-z_t)] d\lambda$$
(1)

where *x* and  $\lambda$  are introduced variables that satisfy  $x_j^2 = \lambda^2 - k_j^2$  and  $k_j^2 = \mu_j \varepsilon_j \omega^2 - i\mu_j \sigma_j \omega$ . In addition,  $\omega$  denotes the angular frequency, and  $I_q(\cdot)$  and  $K_q(\cdot)$  are the first and second modified Bessel functions, respectively, of order *q*. Turning the current off instantaneously causes the magnetic field to decrease suddenly, and an eddy current is then formed on the conductive medium. The electrical parameters, geometric parameters, and boundary conditions of each layer can be used to determine the eddy current field in the *j*th layer:

$$H_2(z_t, \omega, r, z) = \frac{Ir_1}{\pi} \int_0^\infty x_j \left[ -\tau_j K_0(x_j r) I_1(x_j r_1) + C_j I_0(x_j r) - D_j K_0(x_j r) \right] \cos\left[\lambda_j (z - z_t)\right] d\lambda_j$$
(2)

where  $\tau_2 = 1$ ,  $\tau_{j \neq 2} = 0$ , and  $C_j$  and  $D_j$  are coefficients related to the electrical and geometric parameters of each layer of the medium.

According to the *P*-order Gaver–Stehfest inverse Laplace transform formula [24], the magnetic field in the time domain can then be expressed as

$$H(z_t, t, r, z) = \begin{cases} \frac{\ln 2}{t} \sum_{p=1}^{p} E_p H_1(z_t, p \ln 2/it, r, z) & t < t_{\text{off}} \\ \frac{\ln 2}{t} \sum_{p=1}^{p} E_p H_2(z_t, p \ln 2/it, r, z) & t \ge t_{\text{off}} \end{cases}$$
(3)

where  $E_p$  is the Gaver–Stehfest inverse Laplace coefficient, and  $t_{off}$  denotes the moment when the current is turned off. The magnetic field is excited by the transmitting coil when  $t < t_{off}$ , and it produces the eddy current field when  $t > t_{off}$ . The above model represents the application of PEC systems to borehole casing inspection. However, increasing safety concerns in oil and gas production mean that the traditional single-transmitter and singlereceiver probes cannot meet the requirements for challenging downhole environments.

### 3. Synthesized Transmitting Coil

The longitudinal resolution has a significant influence on the performance of PEC systems. One approach to improving the longitudinal resolution is to use MTCs to realize MF [17]. However, MTCs are difficult to fit in a narrow downhole environment because of the complex transmitting array and many electrical circuits.

#### 3.1. Magnetic Focusing

We used the model of the PEC system presented in Section 2 to design the proposed STC for MF in downhole environments. Figure 2 compares the structures of the traditional MTCs and the proposed STC. Both are able to focus a synthesized magnetic field more narrowly on a certain area. As shown in Figure 2a, when MTCs send the same current  $I_1$  simultaneously, a non-focused magnetic field is excited. MTCs can control the vector of the magnetic field (or transmitting currents with  $I_1$ ,  $I_2$ ,  $I_3$ ,  $I_4$ , and  $I_5$ ) to achieve MF. However, MTCs not only require more transmitting electrical circuits, but also require the H-bridges and control signals to be almost the same to ensure consistency when the magnetic field is turned on and off; otherwise, large errors can be introduced. Figure 2b shows the proposed STC, which combines the MTCs into a single coil with multiple sections having different winding directions and number of turns to adjust the magnetic field distribution. Specifically, the STC comprises multiple transmitters, each defined as a section. The five sections are connected together, and each section has a different winding direction and number of turns.



(a)

(b)

Figure 2. Structural comparison: (a) MTCs and (b) the proposed STC.

The sections of the STC can be divided into two types according to the winding direction: the P-sections have the same winding direction as the middle section, and the N-sections have the opposite winding direction of the middle section. If the total number of turns of the STC is *M*, then the magnetic field can be expressed as the sum of the fields of each transmitting coil turn:

$$\overline{H}(\mathbf{z}_{\mathbf{t}}, t, r, z) = \sum_{m=1}^{M} \alpha_m H(z_{t_m}, t, r, z)$$
(4)

where  $z_{tm}$  is the position of the *m*th turn.  $\alpha_m$  is a coefficient related to the winding direction of the *m*th turn coil and can be expressed as:

$$\alpha_m = \begin{cases} 1 & m \text{th turn coil} \in \text{P-Section} \\ -1 & m \text{th turn coil} \in \text{N-Section} \end{cases}$$
(5)

By stacking  $\alpha_m$  to form a vector, we obtain:

$$\boldsymbol{\alpha} = \begin{bmatrix} \alpha_1 & \alpha_2 & \cdots & \alpha_M \end{bmatrix} \tag{6}$$

The winding parameters of the STC can be expressed by  $\alpha$ . Different magnetic field distributions can be achieved by designing the vector  $\alpha$ . Let the number of STC sections be *K*, and the number of turns for each section be  $n_k$ . The two adjacent sections have opposite winding directions. Therefore,  $\alpha$  of the five-section STC shown in Figure 2b can be written as:

$$\boldsymbol{\alpha} = \begin{bmatrix} \underbrace{1 & \cdots & 1}_{n_1} & \underbrace{-1 & \cdots & -1}_{n_2} & \underbrace{1 & \cdots & 1}_{n_3} & \underbrace{-1 & \cdots & -1}_{n_4} & \underbrace{1 & \cdots & 1}_{n_5} \end{bmatrix}$$
(7)

Suppose that  $S_k$  represents the *k*th section for a five-section STC, where  $S_1$ ,  $S_3$ , and  $S_5$  are P-sections, and  $S_2$  and  $S_4$  are N-sections. To focus the magnetic field on the longitudinal center of the probe, the proposed STC adopts a central-symmetry structure. If the number of sections *K* is odd, and symmetrical sections have the same number of turns (e.g.,  $n_1 = n_5$  and  $n_2 = n_4$  for a five-section STC), then the number of turns of section *k* is equal to the number of turns of section K + 1 - k:

$$n_k = n_{K+1-k} \tag{8}$$

Because *K* is an odd number, the central section is the (K + 1)/2-th section. When *K* changes, the sequence number for P-sections also changes. For example,  $S_2$  is a P-section for a three-section STC while  $S_1$ ,  $S_3$ , and  $S_5$  are P-sections for a five-section STC. The P-sections of STCs with different sequence numbers can be expressed as follows:

P-sections = 
$$\begin{cases} S_2, S_4, \dots, S_{(K+1)/2}, \dots, S_{K-1} & K = 3, 7, 11, \dots \\ S_1, S_3, \dots, S_{(K+1)/2}, \dots, S_K & K = 5, 9, 13, \dots \end{cases}$$
(9)

Similar to MTCs, which achieve MF by controlling the magnitude and direction of the current in each element of a multi-coil array, the STC controls the magnetic field distribution according to the winding direction and number of turns in each section. Note that the STC is adjusted by the number of turns in each section. In contrast, the MTC is adjusted by the transmitting current, which offers a finer resolution. However, the STC is a single transmitter and requires a simplified single electrical circuit to realize MF, which is more suitable for harsh downhole environments with limited space.

#### 3.2. Optimization of the STC Parameters Using FoM

In order to optimize the STC parameters, we made a simulation to analyze the MF performance using the STC, MTC, and no-MF. Figure 3 shows the simulation of the longitudinal component of the magnetic field Hz with different winding schemes using the magnetic module of COMSOL 6.0. A standard 5 in casing was used with a wall thickness of 7.52 mm and inner radius r = 55.98 mm, where t = 0 corresponds to the time when the electrical field was turned off. A three-section STC with 300 turns in total was considered. The radius and wire diameter of the transmitting coils were 12 and 0.33 mm, respectively. The second section (S<sub>2</sub>) of the STC was a P-section (with a clockwise winding direction), and S<sub>1</sub> and S<sub>3</sub> were N-sections (with an anticlockwise winding direction) with the same number of turns ( $n_1 = n_3$ ). The proportions for the P- and N-sections of the STC are represented

by gray and black rectangular blocks, respectively. The longitudinal distribution of the magnetic field along the inner wall of the casing was simulated when  $n_2 = 240$  (80% of the total number of turns), 180 (60%), 120 (40%), and 60 (20%). The no-MF and MTC-based MF [17] cases are shown for comparison. The simulation clearly showed that the magnetic field distribution can be changed according to the winding direction of the STC. The MTCs were set to a total number of 300 turns; although the STC had almost the same performance as the MTC-based MF, there should be some performance loss because the MTCs can adjust the transmitting current more finely in theory. To analyze the MF performance, we followed the lead of Tsukada et al. [14] and defined the main lobe width as the distance between the two nearest intersections of the -3 dB line and the curves. Figure 3a shows that the STC obtained a narrower main lobe than the no-MF case when  $n_2 = 240$ . As  $n_2$  was decreased, the main lobe became much narrower, but side lobes appeared, as shown in Figure 3b. At  $n_2 = 120$ , the height of the side lobe exceeded that of the main lobe, as shown in Figure 3c. At  $n_2 = 60$ , the main lobe became wider than that of the no-MF case and was even more deformed, as shown in Figure 3d. Therefore, the proposed STC can obtain a narrower main lobe than the no-MF case, but the appropriate winding parameters need to be determined.



(c)

**Figure 3.** Longitudinal magnetic field distribution at r = 55.98 mm and t = 0 of the three-section STC: (a)  $n_2 = 240$ , (b)  $n_2 = 180$ , (c)  $n_2 = 120$ , and (d)  $n_2 = 60$ .

Similar to Kim et al. [15], we defined an FoM to evaluate the MF performance. We then determined the optimal STC design according to the combination of winding parameters that realized the largest FoM. The MF performance is related to the width of the main lobe and height of the side lobe. The ratios of the side lobe height,  $h_s$ , and reduction in the main lobe width,  $w_r$ , can be defined as follows:

$$h_{\rm s} = \frac{h_{\rm side\ lobe}}{h_{\rm main\ lobe}} \tag{10}$$

$$w_{\rm r} = \frac{w \Big|_{\rm non-focusing} - w \Big|_{\rm focusing}}{w \Big|_{\rm non-focusing}}$$
(11)

where *h* denotes the height of the main lobe or side lobe, and *w* denotes the width of the main lobe in the no-MF, MTC-based MF, or STC-based MF cases. The MF performance is optimized by decreasing  $h_s$  and increasing  $w_r$ . Therefore, the following FoM was defined:

$$FoM = \sqrt{w_r} \times (1 - h_s) \tag{12}$$

where a square root is applied to  $w_r$  to increase its contribution to the FoM because this parameter directly affects the longitudinal distribution of the magnetic field. Note that when the main lobe width is widened (i.e.,  $w_r$  is negative, as shown in Figure 3d), the MF performance is poor. This was outside the scope of this study, and the FoM was set to zero in this case. Thus, the optimal design of the STC can be found by determining the number of sections and total turns according to the winding process and instrument space, as well as a suitable combination of winding parameters.

In order to clarify the design procedure, we can use a three-section STC as an example. Figure 4 shows the values of the FoM,  $w_r$ , and  $h_s$  for different designs using the same simulation parameters shown in Figure 3, where the horizontal axis denotes the variation of the ratio of  $n_2$  to total turns.



Figure 4. MF performance of the proposed STC with different winding parameters.

It can be observed from Figure 4 that decreasing the ratio of  $n_2$  to total turns increased  $w_r$  and  $h_s$ , which decreased the main lobe width and increased the side lobe height. This is consistent with the results shown in Figure 3. In order to optimize the MF performance of the proposed STC, we chose to use the parameters corresponding to the maximum FoM that can be achieved, where FoM = 0.6294 and  $n_2/300 = 0.613$  as shown in Figure 4. As a result, the optimal design parameters of the three-section STC can be acquired, since  $n_1 = n_3$  ( $n_2 = 184$ ,  $n_1 = n_3 = 58$  herein). Similar to the three-section case, the case with more sections can also be extended. It should be noted that the maximum FoM will correspond to different section schemes and the relationship among the coil turns in each section.

# 4. Experiments

For further validation, we performed simulations and experiments using the proposed STC, where the conditions in simulations and experiments were the same. Figure 5 shows the longitudinal magnetic field distribution of the STC using the optimal design in Section 4 at different inner casing radii and times. As shown in Figure 5a, the longitudinal magnetic field distribution of the STC had a narrow main lobe and a short side lobe at {r = 55.98 mm, t = 0 ms}. The magnetic field distribution differed with the time and radii. Figure 5b–d compare the performances of the no-MF case, MTF-based MF, and STC-based MF at {r = 55.98 mm (inner radius of the casing pipe), t = 10 ms}, {r = 59.74 mm (center of wall of the casing pipe), t = 0 ms}, and {r = 63.5 mm (outer radius of the casing pipe), t = 0 ms}, respectively. The designed FoM can be used to select appropriate winding parameters for the STC and optimize the MF performance to approach that of the ideal MTC-based MF. Comparing with the no-MF case, the longitudinal magnetic field distribution at other radii and observation times can also be focused.



(c)

(**d**)

**Figure 5.** Longitudinal magnetic field distributions with no-MF, MTC-based MF, and STC-based MF: (a) r = 55.98 mm, t = 0 ms; (b) r = 55.98 mm, t = 10 ms; (c) r = 59.74 mm, t = 0 ms; and (d) r = 63.5 mm, t = 0 ms.

Without loss of a generality, we further simulated the FoM in a five-section case. Figure 6 shows the MF performance of the five-section STC. The total number of turns was set to 300, and the other parameters were the same as those shown in Figure 4. As shown in Equations (8) and (9),  $S_1$ ,  $S_3$ , and  $S_5$  were P-sections,  $S_2$  and  $S_4$  were N-sections, and  $n_1 = n_5$  and  $n_2 = n_4$ . Compared with the three-section STC in Figure 4, the five-section STC had a more complex design, and its parameters were related to the coil turns of multiple sections; thus, two uncorrelated factors,  $n_3/300$  and  $n_1/(n_1 + n_2)$ , were used to represent the winding parameters of the five-section STC in Figure 6.



Figure 6. Variation in MF performance of the five-section STC with different winding parameters.

From Figure 6, we can conclude that the MF performance was closely related to the probe design. The maximum FoM value of the five-section STC (FoM = 0.7254) was obtained at  $n_3/300 = 0.3867$  and  $n_1/(n_1 + n_2) = 0.2391$  ( $n_3 = 116$ ,  $n_2 = n_4 = 70$ ,  $n_1 = n_5 = 22$ ). The above winding scheme with the optimal FoM was used to design the structures of the three- and five-section STCs, as shown in Figure 7, in which the PEC inspection system used for the experimental prototype was also included.



**Figure 7.** Probe structures and PEC inspection system for the experiments: (**a**) non-focusing probe structure; (**b**) five-section STC probe; (**c**) three-section STC probe; and (**d**) the PEC inspection system.

In Figure 7, the no-MF probe is also shown for comparison. All probes had 300 turns in total. For the three-section STC,  $S_{1-3}$  had 58, 184, and 58 turns, respectively; while for the five-section STC,  $S_{1-5}$  had 22, 70, 116, 70, and 22 turns, respectively. In our experiment, an

observation array with 25 giant magnetoresistance (GMR) sensors (AAH002-02 from NVE Corporation, Minneapolis, MN, USA) was used to measure the longitudinal magnetic field on the inner wall of the casing.

Figure 8 shows the normalized magnetic field of the no-MF, MTC-based MF, and STC-based MF cases with three and five sections in the experiment and simulation. The experimental and simulation conditions were identical.



Figure 8. Experimental and simulation results for MF with the five-section STC.

As shown in Figure 8, the experimental results for the no-MF, MTC-based MF, and STC-based MF cases showed good agreement with the simulation results. The error can be attributed to the inductance error of the additional inductor [15] and differences in the coil geometry. Additionally, the main lobe of the STC was significantly narrower than the no-MF and approximate to that of the MTC for both the three-section and five-section STCs; as the number of sections in STC increases, the main lobe becomes much narrower. Although the proposed STC can realize MF by narrowing the main lobe, the overall magnetic field intensity is reduced because the N-sections excite the magnetic field in the opposite direction. In other words, the STC is a tradeoff that sacrifices magnetic field intensity for a narrower main lobe.

To clarify the tradeoff relationship, we can use the peak value of the main lobe  $p_m$  to characterize the total magnetic field intensity. Figure 9 compares the maximum FoMs that can be achieved by STC- and MTC-based probes and their corresponding  $w_r$  values. The maximum current of the transmitting coil was set to 1 A, and the total number of turns was 300 for all probes. The three- and five-section STCs were designed according to the optimal parameters given in Figures 4 and 6, respectively.

From Figure 9, it can be observed that increasing  $p_{\rm m}$  clearly decreased the MF effect with respect to FoM and  $w_{\rm r}$ . Note that  $p_{\rm m}$  was a constant value to ensure that the eddy current field was sufficiently large enough to meet detection range and SNR requirements. Setting  $p_{\rm m}$  higher is not required because the casing is not very large, and setting  $p_{\rm m}$  lower would not result in an acceptable SNR. Specifically, when  $p_{\rm m} = 0$ , the STC- and MTC-based probes achieved their maximum FoM values regardless of the magnetic field intensity. The FoM and  $w_{\rm r}$  values of both STCs nearly reached that of the MTC. In addition, the FoM and  $w_{\rm r}$  values of both the three- and five-section STCs remained unchanged, with increasing  $p_{\rm m}$  until  $p_{\rm m} = 0.289 \times 10^{-3}$  T and  $1.592 \times 10^{-3}$  T, respectively. The five-section STC had a higher FoM than the three-section STC at the cost of a reduced magnetic field intensity. However, when  $p_{\rm m} > 1.592 \times 10^{-3}$  T, the two STCs had very similar FoMs because most of the coil turns were allocated to the middle section. In most cases shown in Figure 9, the FoMs of the STC nearly reached that of the MTC-based probe. At  $p_{\rm m} = 2.75 \times 10^{-3}$  T,

the FoMs of the STCs and MTC-based probes were equal because the N-section had zero turns to maximize  $p_m$ . In general, the MF performance of STC with a single transmitter and a simplified single electrical circuit was close to that of the MTC, which can reach the needed sensitivity for the detection of submillimeter defects, and is more suitable for harsh downhole environments with limited space.



**Figure 9.** FoMs of the MTC and STC at different  $p_{\rm m}$  values.

## 5. Conclusions

A novel STC is proposed comprising connected sections with different winding directions and number of turns to realize a PEC system with MF for downhole casing inspection. The proposed design simplifies the electrical circuit and is more suitable for narrow downhole environments. A model of the PEC system was used to theoretically derive the transmitting magnetic field, and an FoM was constructed to evaluate the MF performance and optimize the winding parameters of the STC. The simulation and experimental results for a standardized oil well casing revealed that the proposed STC is effective at realizing MF.

Author Contributions: Conceptualization, C.L.; methodology, B.D.; software, C.L.; validation, Y.Z.; formal analysis, J.D.; investigation, L.Y.; resources, B.D.; data curation, L.Y.; writing—original draft preparation, L.Y.; writing—review and editing, L.Y.; visualization, R.D.; supervision, Y.Z.; project administration, R.D.; funding acquisition, B.D. All authors have read and agreed to the published version of the manuscript.

**Funding:** This work was supported in part by the National Natural Science Foundation of China under Grant 51974250 and Grant 41874158, in part by the Youth Science and Technology Nova Project in Shaanxi Province, China, under Grant 2020KJXX-018, and in part by the Postgraduate Innovation and Practical Ability Training Program of Xi'an Shiyou University under Grant YCS22111001.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

**Data Availability Statement:** All the details of this work, including data and algorithm codes, are available by contacting the corresponding author: dangbo@xsyu.edu.cn.

**Acknowledgments:** The authors would like to thank the reviewers for their helpful suggestions, which have considerably improved the quality of the manuscript.

Conflicts of Interest: The authors declare no conflict of interest.

## References

- 1. Malekmohammadi, H.; Migali, A.; Laureti, S.; Ricci, M. A Pulsed Eddy Current Testing Sensor Made of Low-Cost Off-the-Shelf Components: Overview and Application to Pseudo-Noise Excitation. *IEEE Sens. J.* **2021**, *21*, 23578–23587. [CrossRef]
- Meng, T.; Tao, Y.; Chen, Z.; Avila, J.R.S.; Ran, Q.; Shao, Y.; Huang, R.; Xie, Y.; Zhao, Q.; Zhang, Z.; et al. Depth Evaluation for Metal Surface Defects by Eddy Current Testing Using Deep Residual Convolutional Neural Networks. *IEEE Trans. Instrum. Meas.* 2021, 70, 1–13. [CrossRef]
- 3. Yan, B.; Li, Y.; Liu, Z.; Ren, S.; Chen, Z.; Lu, X.; Abidin, I.M.Z. Pulse-Modulation Eddy Current Imaging for 3D Profile Reconstruction of Subsurface Corrosion in Metallic Structures of Aviation. *IEEE Sens. J.* 2021, *21*, 28087–28096. [CrossRef]
- Niu, C.; Wang, L.; Wang, Y.; Zhai, Y.; Qu, H.; Liu, Y.; Wang, Q. Numerical Analysis of Eddy Current Induced by z-Gradient Coil in a Superconducting MRI Magnet. *IEEE Trans. Appl. Supercond.* 2020, 30, 1–6. [CrossRef]
- 5. Xiao, Q.; Feng, J.; Xu, Z.; Zhang, H. Receiver Signal Analysis on Geometry and Excitation Parameters of Remote Field Eddy Current Probe. *IEEE Trans. Ind. Electron.* 2021, *69*, 3088–3098. [CrossRef]
- 6. Suresh, V.; Abudhahir, A.; Daniel, J. Characterization of Defects on Ferromagnetic Tubes Using Magnetic Flux Leakage. *IEEE Trans. Magn.* 2019, 55, 1–10. [CrossRef]
- Wang, Y.; Nie, Y.; Qi, P.; Zhang, N.; Ye, C. Inspection of Defect Under Thick Insulation Based on Magnetic Imaging with TMR Array Sensors. *IEEE Trans. Magn.* 2021, 58, 1–10. [CrossRef]
- Liu, C.; Dang, B.; Wang, H.; Shen, X.; Yang, L.; Ren, Z.; Dang, R.; Kang, Y.; Sun, B. Multiple-Transmit Focusing for the Nondestructive Testing of Downhole Casings Based on Borehole Transient Electromagnetic Systems. *IEEE Access* 2020, *8*, 210978–210987. [CrossRef]
- Fu, Y.; Yu, R.; Peng, X.; Ren, S. Investigation of casing inspection through tubing with pulsed eddy current. *Nondestruct. Test. Eval.* 2012, 27, 353–374. [CrossRef]
- 10. Wu, T.; Bowler, J.; Theodoulidis, T.P.; Th, T. Eddy Current Induction by a Coil Whose Axis is Perpendicular to that of a Tube. *IEEE Trans. Magn.* **2017**, *53*, 6201709. [CrossRef]
- 11. Zhou, H.; Hou, K.; Pan, H.; Chen, J.; Wang, Q. Study on the Optimization of Eddy Current Testing Coil and the Defect Detection Sensitivity. *Press. Vessel. Technol. Prep. Future* **2015**, *130*, 1649–1657. [CrossRef]
- 12. Dang, B.; Yang, L.; Du, N.; Liu, C.; Dang, R.; Wang, B.; Xie, Y. Auxiliary Sensor-Based Borehole Transient Electromagnetic System for the Nondestructive Inspection of Multipipe Strings. *Sensors* **2017**, *17*, 1836. [CrossRef] [PubMed]
- 13. Yang, C.; Gao, B.; Ma, Q.; Xie, L.; Tian, G.Y.; Yin, Y. Multi-Layer Magnetic Focusing Sensor Structure for Pulsed Remote Field Eddy Current. *IEEE Sens. J.* 2018, 19, 2490–2499. [CrossRef]
- Tsukada, K.; Hirata, T.; Goda, Y.; Sakai, K.; Kiwa, T. Hybrid Magnetic Sensor Combined With a Tunnel Magnetoresistive Sensor and High-Temperature Superconducting Magnetic-Field-Focusing Plates. *IEEE Trans. Appl. Supercond.* 2018, 29, 8800405. [CrossRef]
- 15. Kim, J.H.; Choi, B.; Kim, H.R.; Rim, C.-T.; Kim, Y.-S. Single-Variable-Input Active Sidelobe Suppression Method for Synthesized Magnetic Field Focusing Technology and Its Optimization. *IEEE Trans. Ind. Electron.* **2019**, *67*, 9813–9823. [CrossRef]
- 16. Kim, J.H.; Choi, B.H.; Kim, H.R.; Rim, C.T. 2-D synthesized magnetic field focusing technology With loop coils distributed in a rectangular formation. *IEEE Trans. Ind. Electron.* **2019**, *66*, 5558–5566. [CrossRef]
- 17. Liu, C.; Dang, B.; Wang, H.Y.; Yang, L.; Dang, J.; Shen, X.; Dang, R. Synthesised magnetic field focusing for the non-destructive testing of oil and gas well casing pipes using pulsed eddy-current array. *IEEE Trans. Magn.* **2022**. [CrossRef]
- Cubero, D.; Marmugi, L.; Renzoni, F. Exploring the limits of magnetic field focusing: Simple planar geometries. *Results Phys.* 2020, 19, 103562. [CrossRef]
- 19. Tanaka, H.; Iizuka, H. Kilohertz Magnetic Field Focusing Behavior of a Single-Defect Loop Array Characterized by Curl of the Current Distribution With Delta Function. *IEEE Antennas Wirel. Propag. Lett.* **2012**, *11*, 1088–1091. [CrossRef]
- Mohammed, A.I.; Oyeneyin, B.; Atchison, B.; Njuguna, J. Casing structural integrity and failure modes in a range of well types—A review. J. Nat. Gas Sci. Eng. 2019, 68, 102898. [CrossRef]
- Wu, S.N.; Zhang, L.B.; Fan, J.C.; Zhang, X.M.; Zhou, Y.F.; Wang, D.G. Prediction analysis of downhole tubing leakage location for offshore gas production wells (in English). *Measurement* 2018, 127, 546–553. [CrossRef]
- 22. Ulapane, N.; Thiyagarajan, K.; Miro, J.V.; Kodagoda, S. Surface Representation of Pulsed Eddy Current Sensor Signals for Improved Ferromagnetic Material Thickness Quantification. *IEEE Sens. J.* **2020**, *21*, 5413–5422. [CrossRef]
- 23. Kaufman, A.A.; Yu, A. Dashevsky. In Principles of Induction Logging; Elsevier: Amsterdam, The Netherlands, 2003.
- 24. Li, J.H.; Farquharson, C.G.; Hu, X.Y. Three effective inverse Laplace transform algorithms for computing time domain electromagnetic responses (in English). *Geophysics* 2016, *81*, E113–E128. [CrossRef]